



**CHALMERS**  
UNIVERSITY OF TECHNOLOGY

*Alcator  
C-Mod*

## Modeling synchrotron radiation images of runaway electrons

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**O. Jakobsson<sup>1</sup>, T. Fülöp<sup>1</sup>**

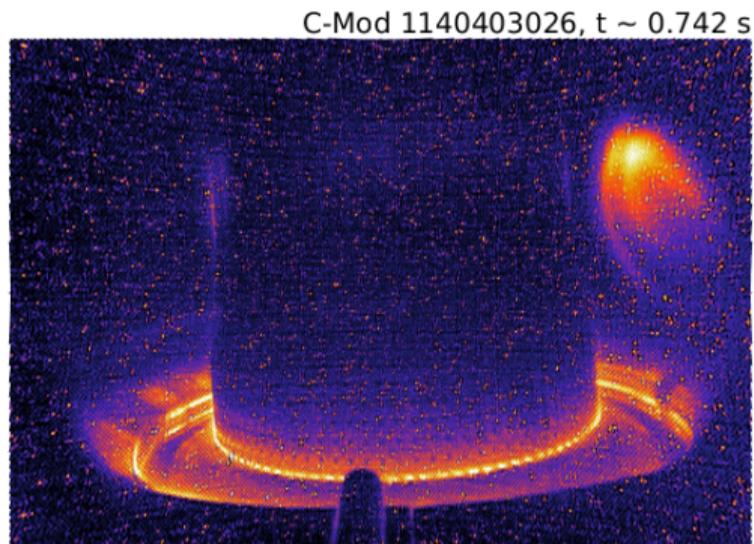
<sup>1</sup> Chalmers University of Technology, Gothenburg, Sweden

<sup>2</sup> PSFC, Massachusetts Institute of Technology, Cambridge, Massachusetts, USA

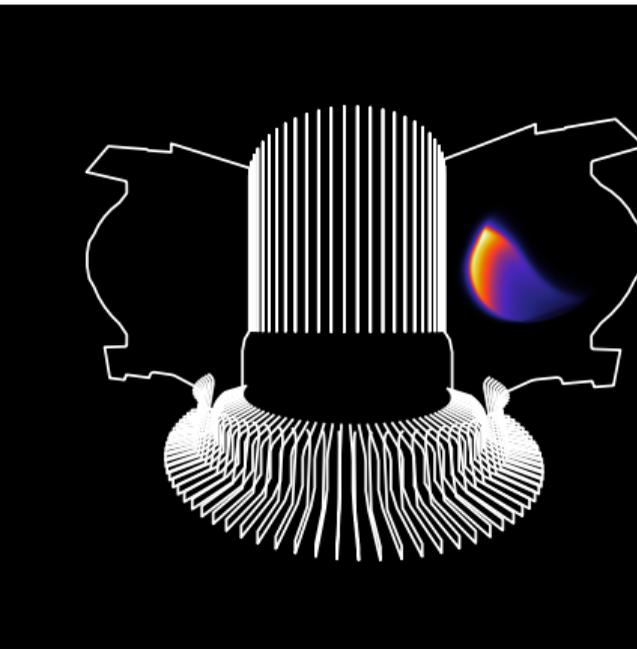


**Q:** *Can we measure the runaway electron distribution function?*

1. SOFT
2. Synchrotron radiation
3. Interpretation of synchrotron spot
4. Case studies
  - I. DIII-D
  - II. Alcator C-Mod



*Synchrotron spot in Alcator C-Mod*



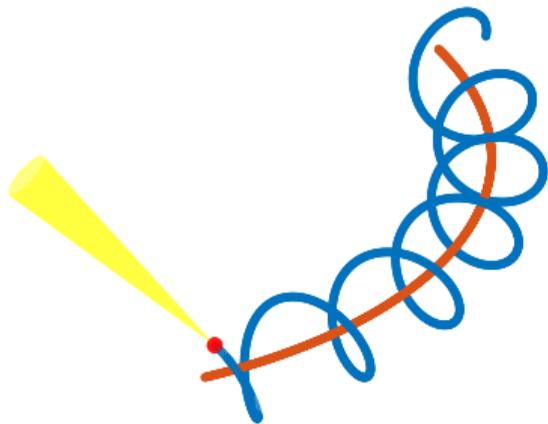
Available on GitHub:  
<https://github.com/hoppe93/SOFT>

- **Sparking question:** How do orbits affect the synchrotron radiation?
- **Particle simulations:**
  - ▶ Drop particles in outer midplane
  - ▶ Follow guiding-center orbits
  - ▶ Register contributions to a detector
- **Main differences to previous models:**
  - ▶ Utilizes numeric magnetic field
  - ▶ Considers *detected* radiation (not just *emitted* [1]) and *intensity variations* (not just *shape* [2,3])

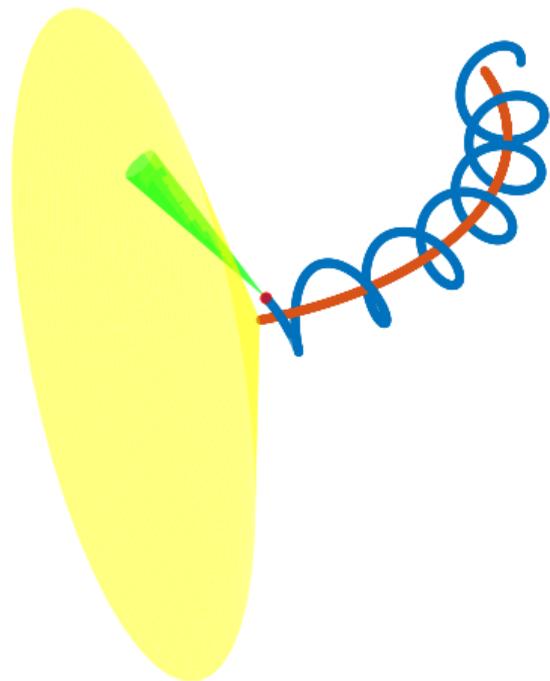
[1] A. Stahl et al., Phys. Plasmas 20 (9), 093302 (2013).

[2] I. M. Pankratov, Plasma Phys. Rep. 22 (6), 535-538 (1996).

[3] R. J. Zhou et al., Phys. Plasmas 21 (6), 063302 (2014).



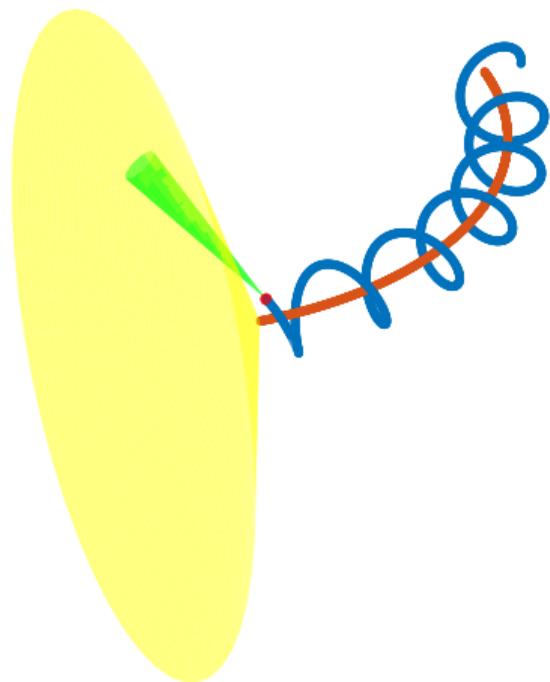
- From highly relativistic particles
- Emitted mainly in the direction of motion
- Angular spread around velocity vector  $\sim \gamma^{-1}$
- Peaks in visible - IR wavelength ranges



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⇒ **Only electrons moving towards observer are visible**

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**Reminder:**

- Assume: guiding-center follows magnetic field lines
- Cone of radiation around GC; opening angle  $\theta_p$

Magnetic field  
unit vector

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$$\frac{\mathbf{x}_0 - \mathbf{x}}{|\mathbf{x}_0 - \mathbf{x}|}$$

$\mathbf{x}_0 =$  Observer location

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The diagram shows the equation  $|\hat{\mathbf{b}}(\mathbf{x}) \cdot \hat{\mathbf{r}}| = |\cos \theta_p|$  with several annotations. A red line points from the text 'Magnetic field unit vector' to the  $\hat{\mathbf{b}}(\mathbf{x})$  term. A blue line points from the text  $\frac{\mathbf{x}_0 - \mathbf{x}}{|\mathbf{x}_0 - \mathbf{x}|}$  to the  $\hat{\mathbf{r}}$  term. A cyan line points from the text 'Point of interest / particle position' to the  $\mathbf{x}$  argument of  $\hat{\mathbf{b}}$ . A yellow circle highlights the  $\theta_p$  term, with a yellow line pointing to the text 'Pitch angle' and the equation  $\cos \theta_p = v_{\parallel} / v$ .

$$|\hat{\mathbf{b}}(\mathbf{x}) \cdot \hat{\mathbf{r}}| = |\cos \theta_p|$$

Magnetic field unit vector

Point of interest / particle position

Pitch angle  
 $\cos \theta_p = v_{\parallel} / v$

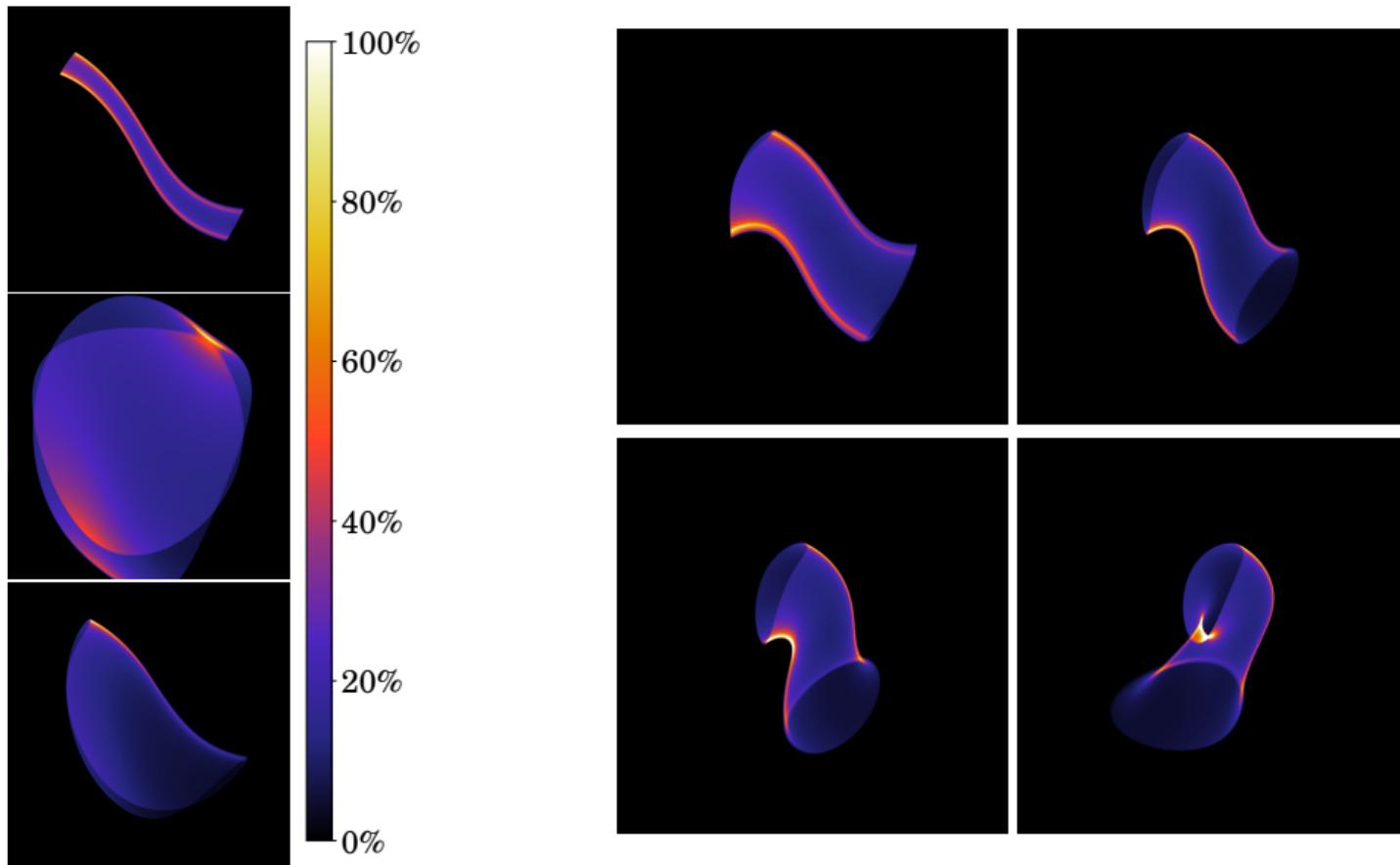
$\frac{\mathbf{x}_0 - \mathbf{x}}{|\mathbf{x}_0 - \mathbf{x}|}$

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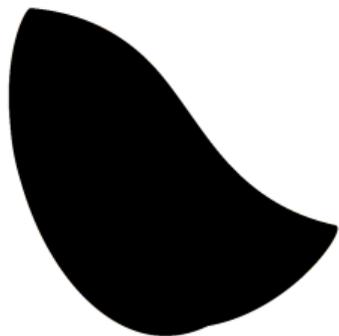
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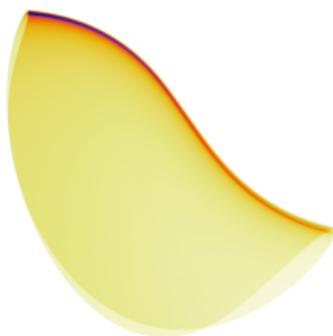
# Surface-of-visibility



**Surface-of-visibility**



**Line-integration  
and emission  
effects**

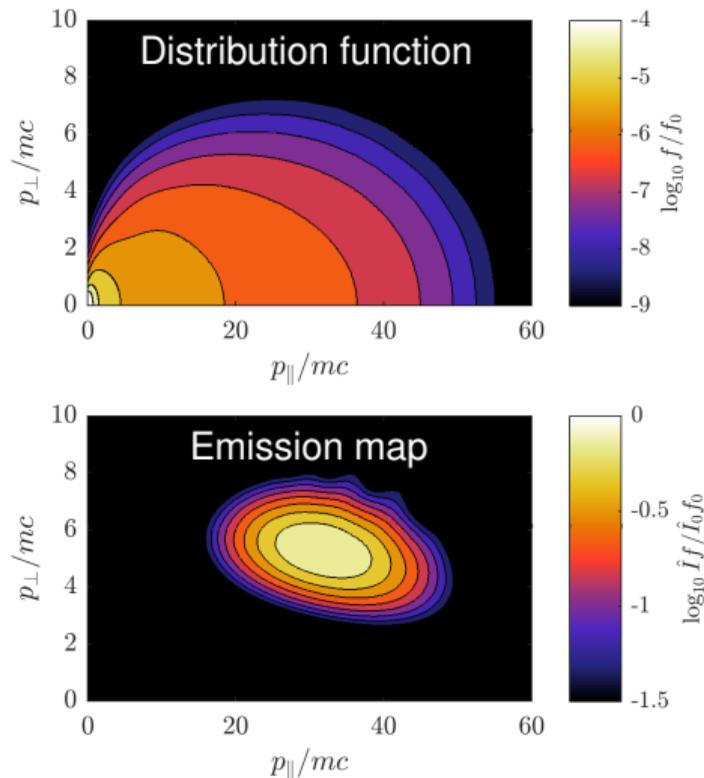


**Finite spectral  
range effects**



Simple spot shape model could have been okay, except...

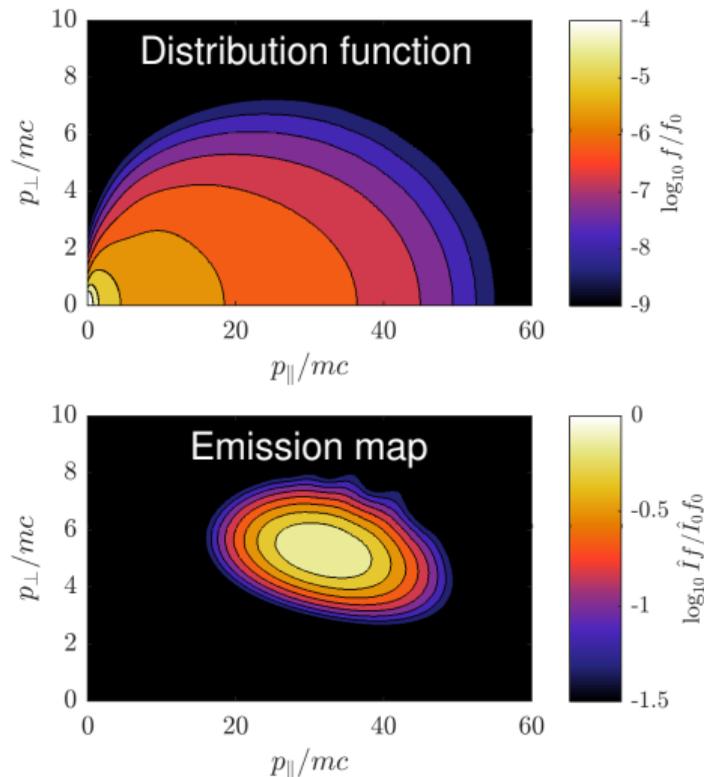
- Particles are widely distributed in momentum-space
- Actual spot shape is a superposition of images



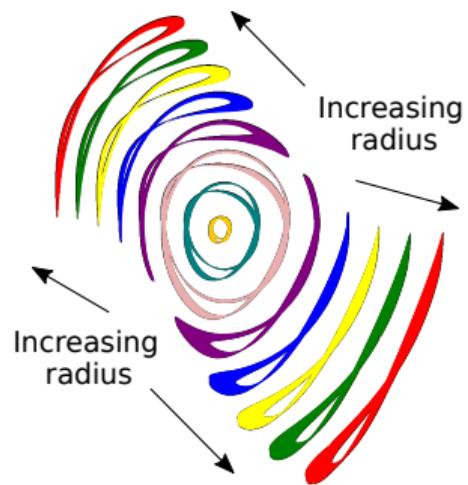
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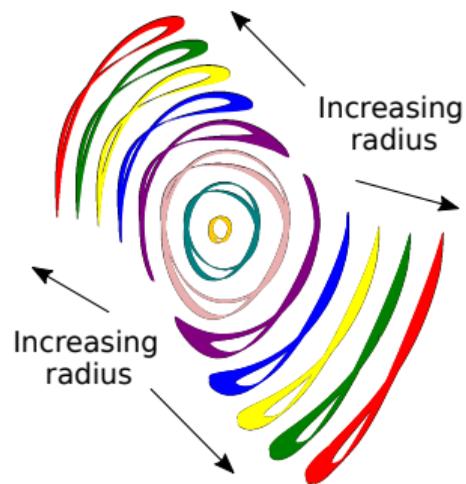
⇒ Image dominated by line-integration and spectral effects



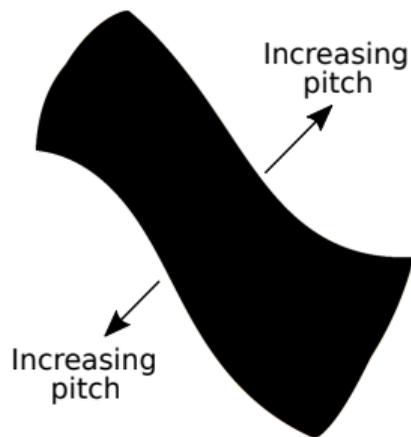
## Radius



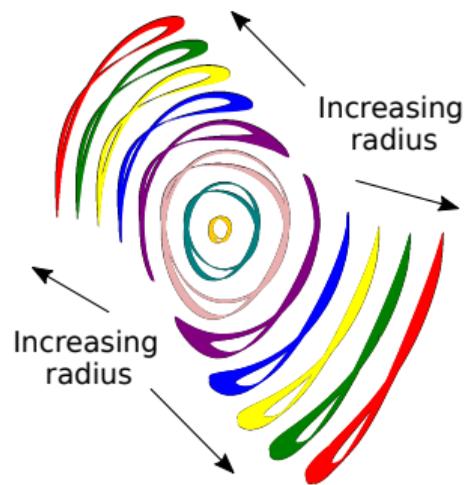
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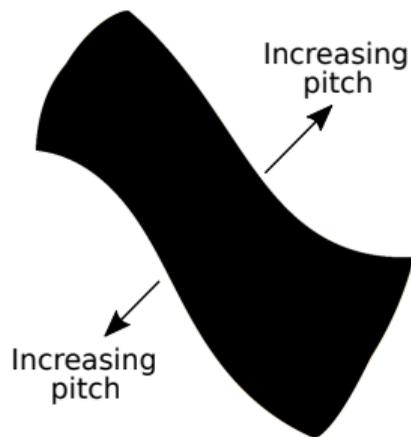
## Pitch angle



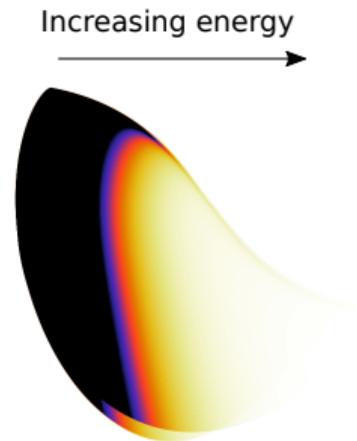
## Radius



## Pitch angle



## Energy



Assume

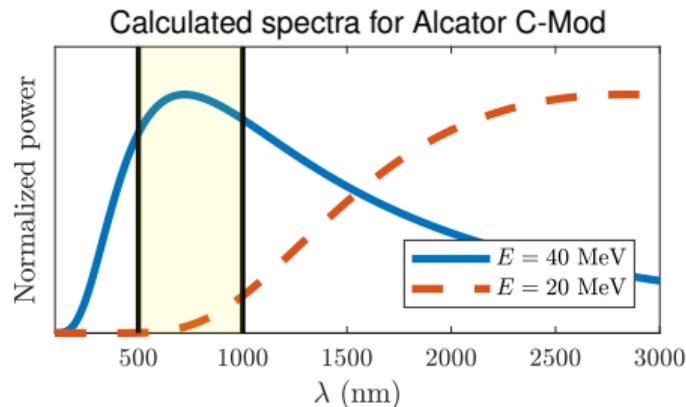
$$B = \frac{B_0 R_0}{R}.$$

In the limit  $\lambda \ll \lambda_c$  the emitted power scales with major radius  $R$  as

$$P \sim \left(\frac{R}{R_c}\right)^{3/4} \exp\left[-\left(\frac{R}{R_c}\right)^{3/2}\right],$$

$$R_c \equiv \frac{B_0 R_0}{2m} \left(\frac{3e\gamma\lambda\sqrt{\mu}}{\pi c^2}\right)^{2/3}.$$

As  $\lambda \rightarrow \lambda_c$ , the scaling gets weaker.

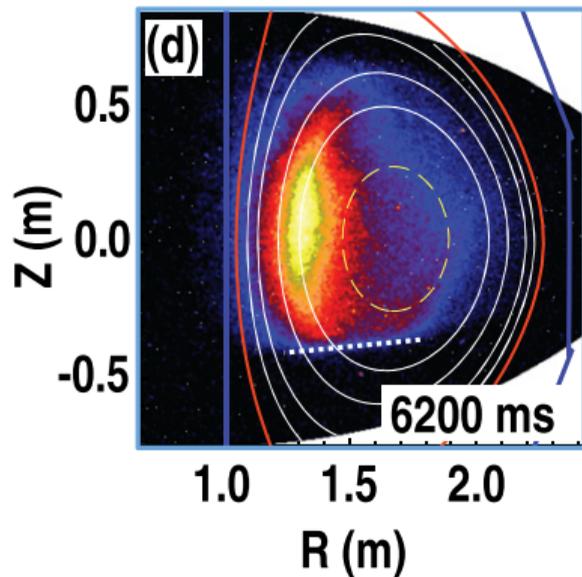


**Alcator C-Mod**

$$R_c \sim \frac{1}{5} R_0$$

**DIII-D**

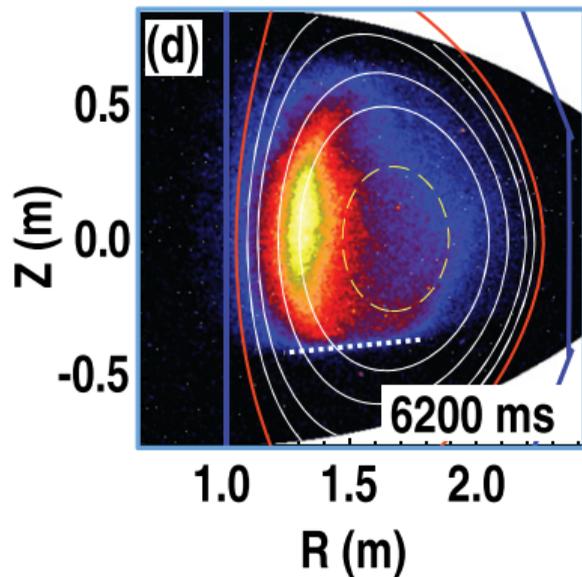
$$R_c \sim \frac{1}{10} R_0$$



## Observations

- Strong radiation from HFS
- Large vertical extent of spot

Experimental image from *C. Paz-Soldan et al., Phys. Plasmas 21 (2), 022514 (2014)*.



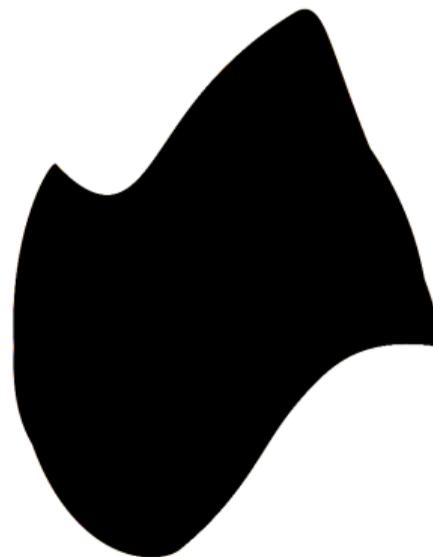
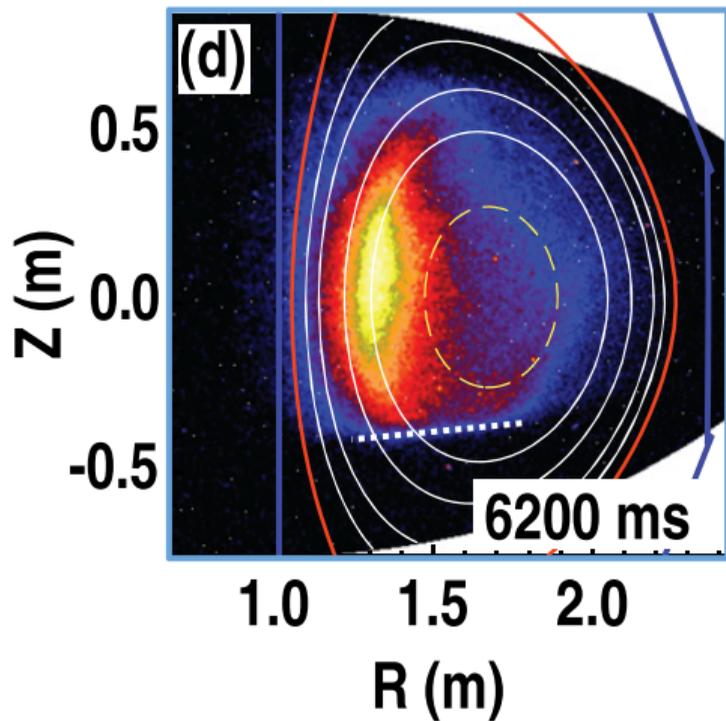
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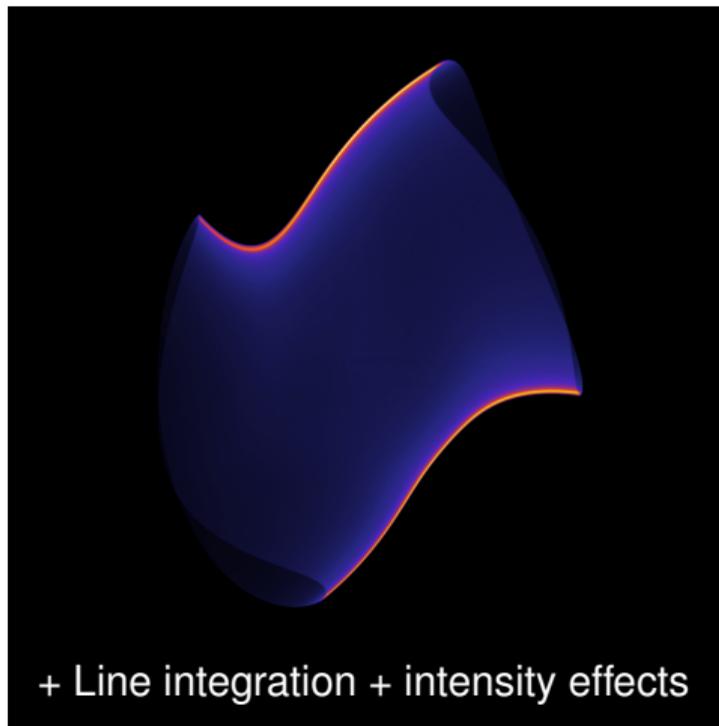
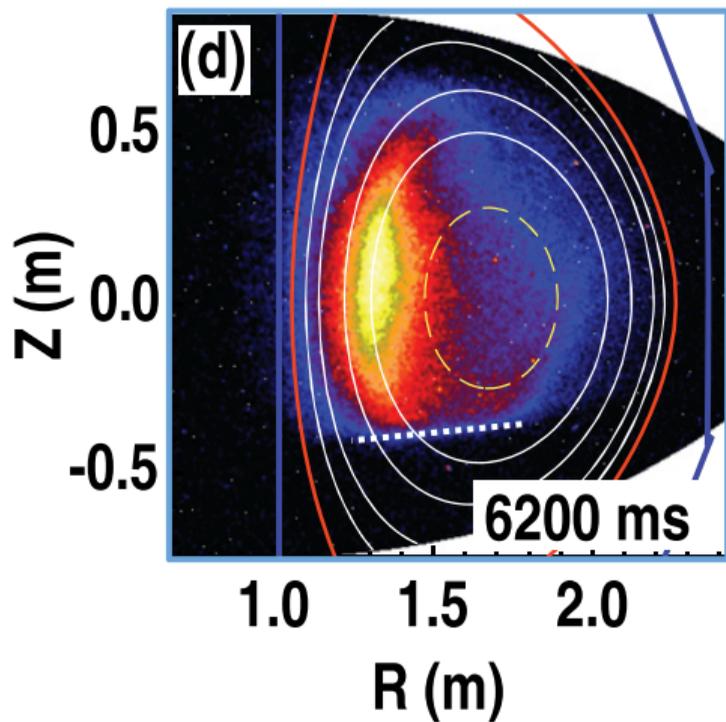
- Dominant energy between 10-20 MeV
- Dominant pitch angles at least 0.20 rad

Experimental image from *C. Paz-Soldan et al., Phys. Plasmas 21 (2), 022514 (2014)*.

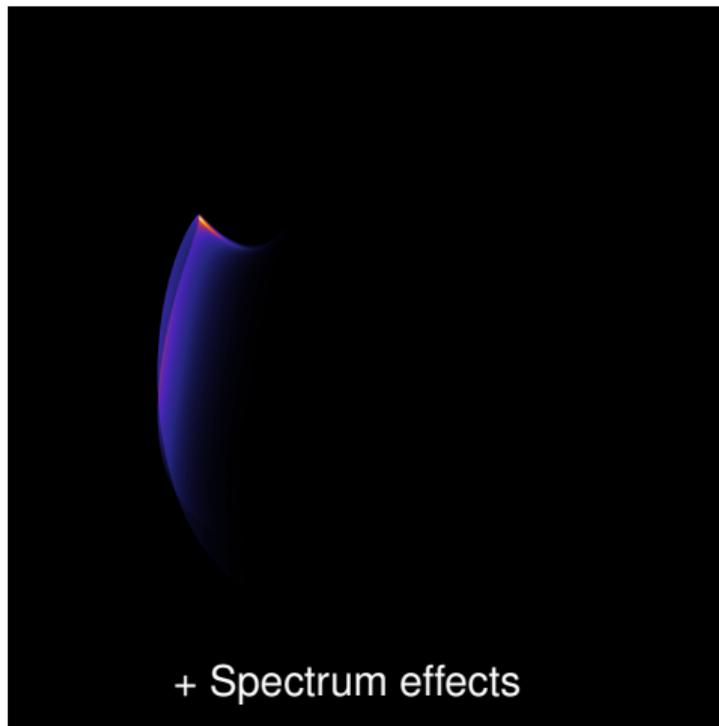
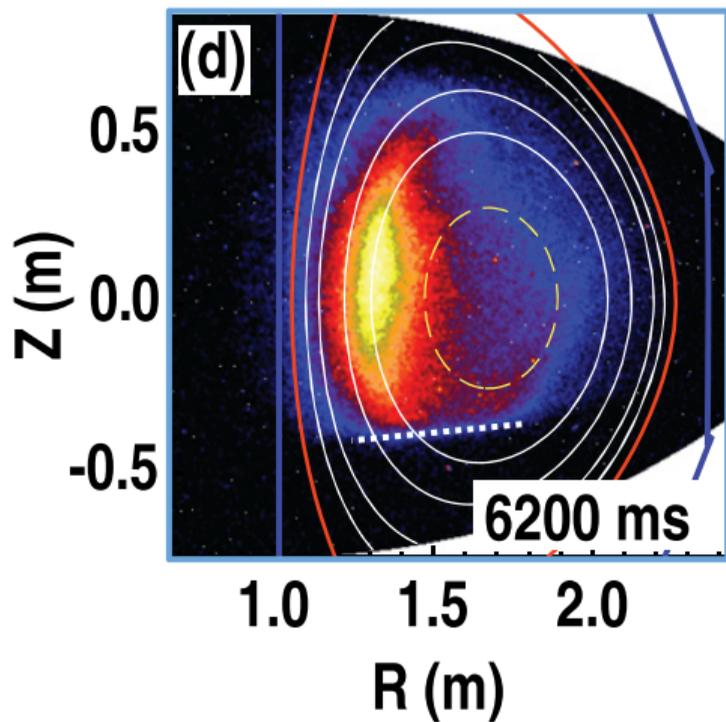


Spot shape

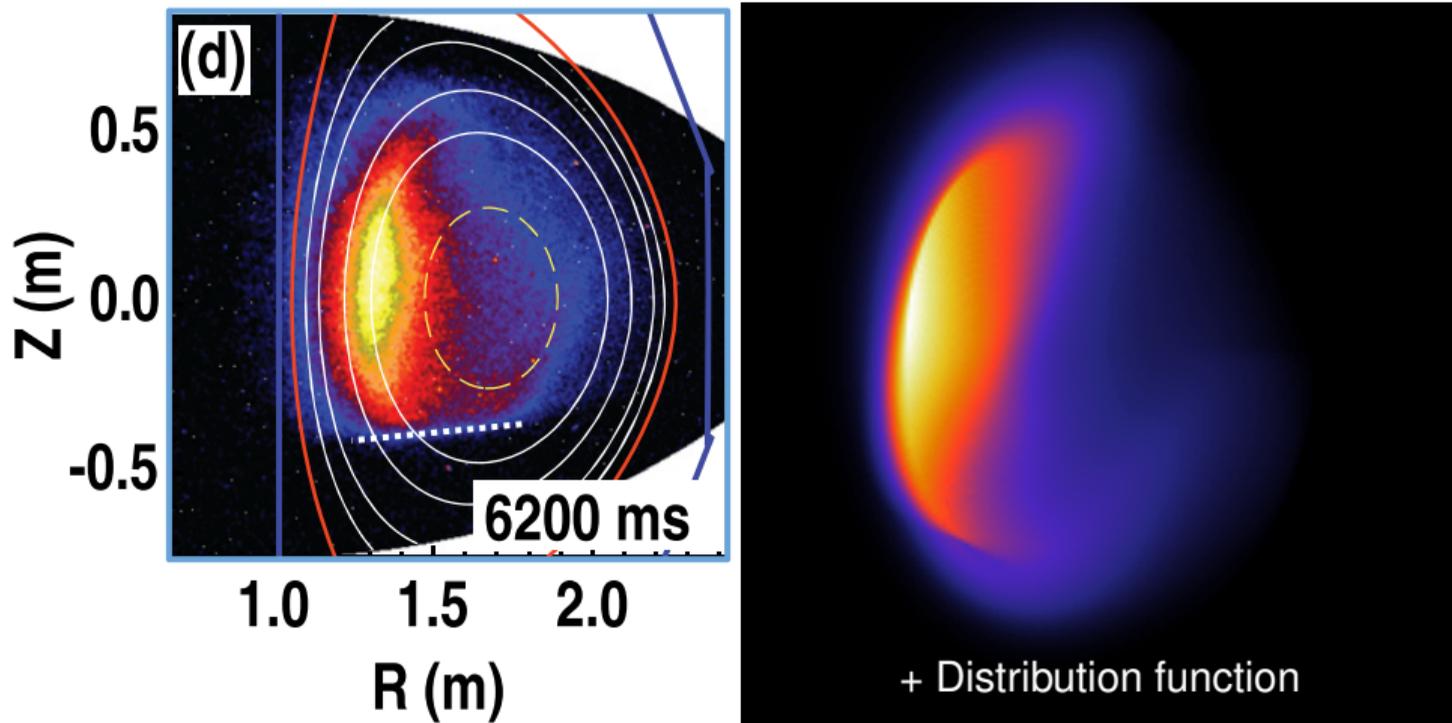
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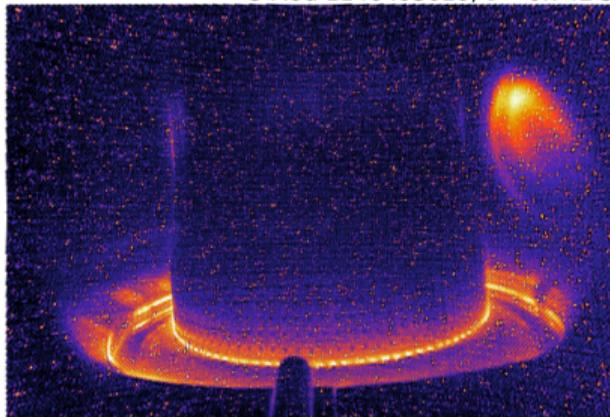


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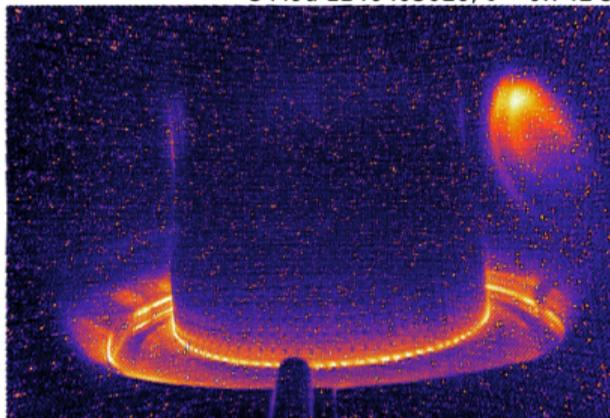
C-Mod 1140403026,  $t \sim 0.742$  s



### Observations

- Horizontally centered radiation
- Three distinct “legs”
- Bright spot at the top

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- Weak spot curvature on the left  $\implies$  dominant pitch angles between 0.10-0.15 rad

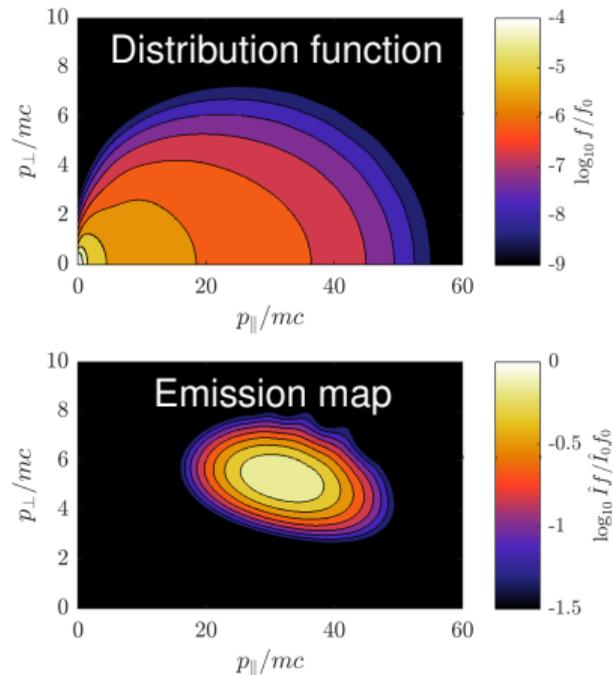
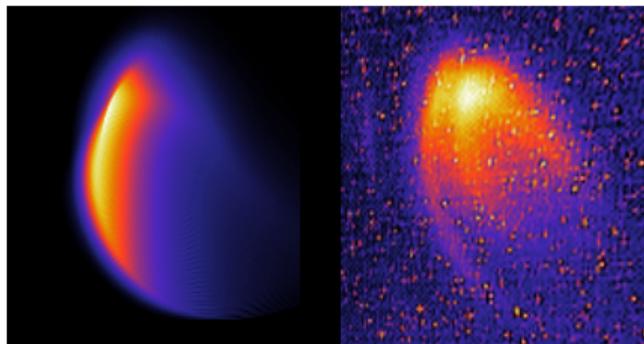
The diagram shows the Fokker-Planck equation for relativistic electrons with several terms highlighted and labeled:

- $\frac{\partial f}{\partial t}$ : Distribution function (purple label)
- $-\frac{e\mathbf{E}}{m_e c} \cdot \frac{\partial f}{\partial \mathbf{p}}$ : Synchrotron losses (orange label)
- $+\frac{\partial}{\partial \mathbf{p}} \cdot (\mathbf{F}_{\text{syn}} f)$ : Fokker-Planck operator (blue label)
- $= C\{f\}$ : Runaway avalanche source (green label)
- $+ S_{\text{ava}}$ : Runaway avalanche source (green label)
- $+ S_h$ : Heat sink (red label)

- Solves Fokker-Planck equation for relativistic electrons
- 0D in space, 2D in momentum space

[4] M. Landreman et al., *Comp. Phys. Comm.* 185 (3), 847-855 (2014).

[5] A. Stahl et al., *Nucl. Fusion* 56 (11), 112009 (2016).

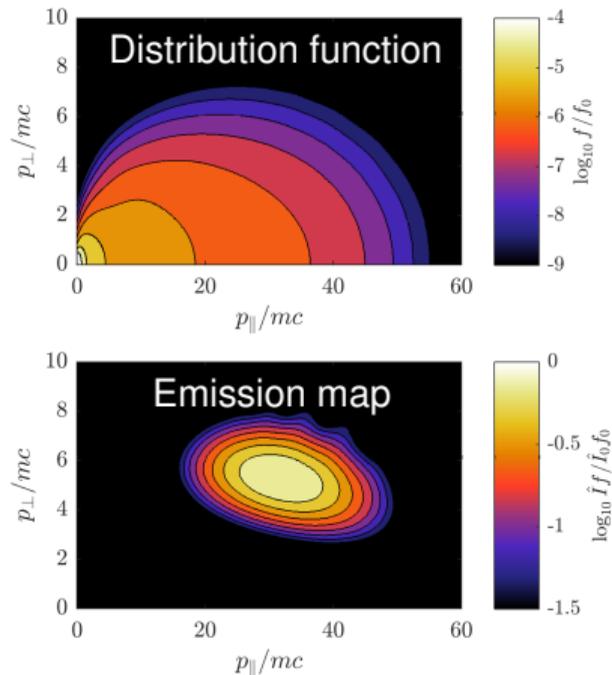
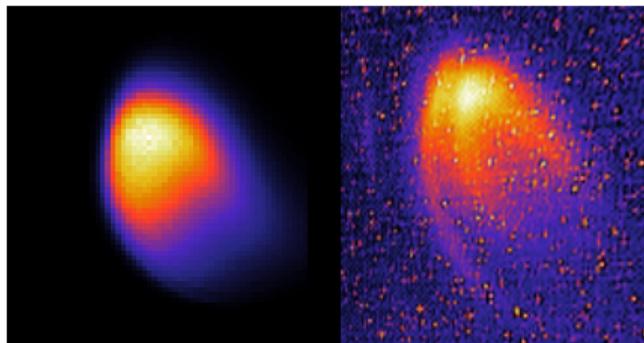


## From kinetic simulations:

- Dominant energy around 15 MeV
- Dominant pitch angle  $\theta_p \sim 0.15$  rad

## Deduced from image:

- Decreasing radial profile
- Sudden and brief “well” in outer **radial profile**

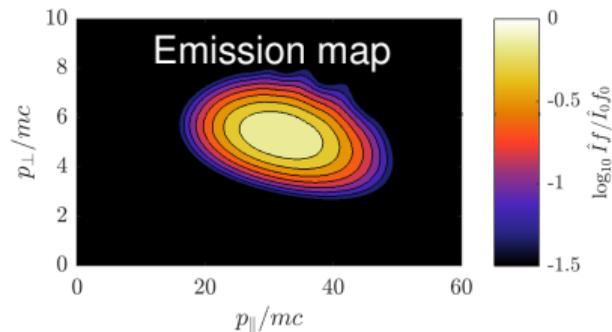
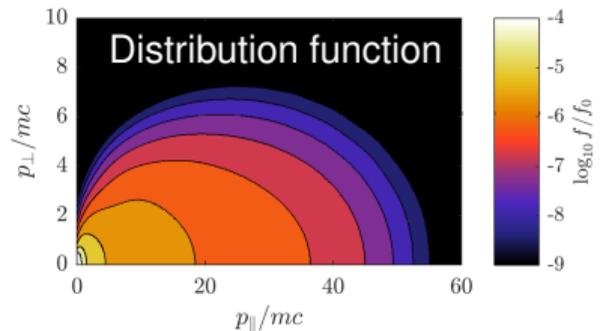
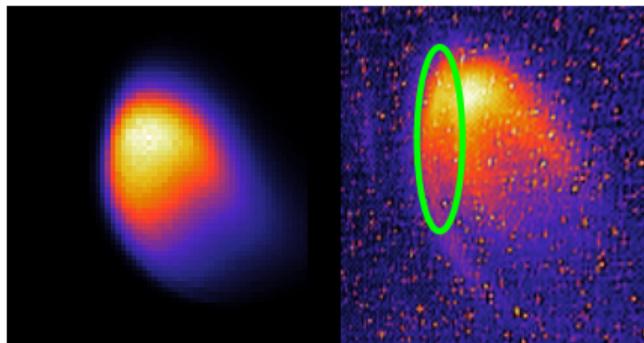


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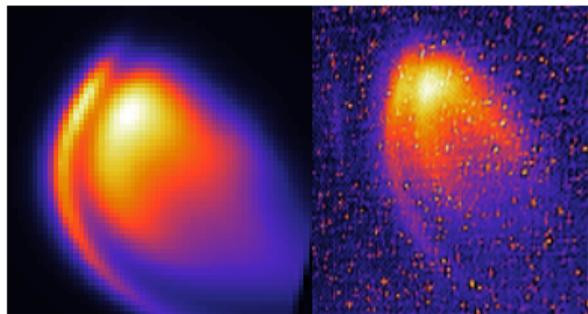
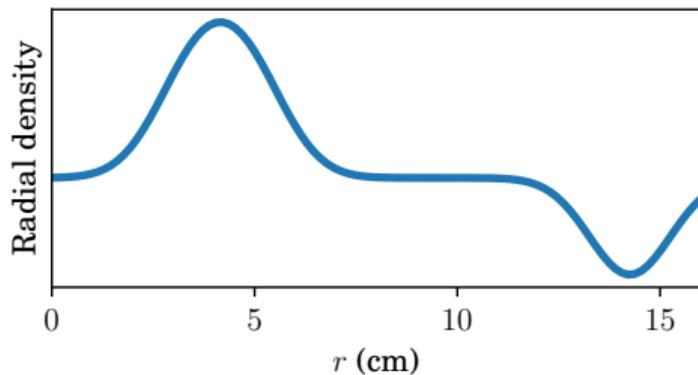
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- Attempt at inverting radial distribution function
- Image comparison + stochastic optimization



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## Further reading

[6] M. Hoppe et al., *SOFT: A synthetic synchrotron diagnostic for runaway electrons*, submitted to Nucl. Fusion, arXiv:1709.00674.